
Modern approaches to quantum gravity

Homework 11

Fall 2025

1. Scalar field in AdS

Consider Minkowskian AdS_{d+1} in conformal coordinates

$$ds^2 = \frac{1}{\cos^2 \theta} (-dt^2 + d\theta^2 + \sin^2 \theta d\Omega_{d-1}^2) \quad (1)$$

We will find the spectrum of a scalar particle whose field obeys the Klein-Gordon equation

$$\square \phi = m^2 \phi \quad (2)$$

where \square is taken in AdS_{d+1} .

- (a) Argue, from the isometries of AdS_{d+1} , that a general solution of the above equation is given by a linear combination of the solutions

$$\phi = e^{-i\omega t} G(\theta) Y_\ell(\Omega) \quad (3)$$

where Y_ℓ are the spherical Harmonics on the $d-1$ sphere S^{d-1} , with coordinates Ω , and $G(\theta)$ is for now an arbitrary function of θ .

- (b) Show that with the above Ansatz, the Klein-Gordon equation reduces to

$$\omega^2 \cos^2 \theta G(\theta) + \frac{d-1}{\tan \theta} G'(\theta) + \cos^2 \theta G''(\theta) - \frac{\ell(\ell+d-2)}{\tan^2 \theta} G(\theta) - m^2 G(\theta) = 0 \quad (4)$$

- (c) Show that the above equation are solved by the solutions G_\pm , where

$$G_+ = (\sin \theta)^\ell (\cos \theta)^\Delta {}_2F_1\left(\frac{\Delta + \ell + \omega}{2}, \frac{\Delta + \ell - \omega}{2}; \ell + \frac{d}{2}; \sin^2 \theta\right) \quad (5)$$

$$G_- = (\sin \theta)^{(2-d-\ell)} (\cos \theta)^\Delta {}_2F_1\left(\frac{\Delta + 2 - d - \ell + \omega}{2}, \frac{\Delta + 2 - d - \ell - \omega}{2}; 2 - \ell - \frac{d}{2}; \sin^2 \theta\right) \quad (6)$$

$$\Delta \equiv \frac{d}{2} + \frac{1}{2} \sqrt{d^2 + 4m^2} \quad (7)$$

Hint: Use Mathematica.

Argue that asking the regularity at $\theta = 0$, $G'(\theta = 0) = 0$ discards G_- .

- (d) Consider the stress-energy tensor

$$T_{\mu\nu} = 2\partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} ((\partial\phi)^2 + m^2 \phi^2) + \beta (g_{\mu\nu} \square - \nabla_\mu \nabla_\nu + R_{\mu\nu}) \phi^2 \quad (8)$$

We require that the energy-momentum flux through the boundary ($\theta = \pi/2$) vanishes

$$\int_{S^{d-1}} d\Omega_{d-1} \sqrt{g} n_i T_0^i \Big|_{\theta=\pi/2} = 0 \quad (9)$$

Show that this requirement gives the boundary condition

$$(\tan \theta)^p [(1 - 2\beta)\partial_\theta + 2\beta \tan \theta]\phi^2 \rightarrow 0 \quad (\theta \rightarrow \pi/2) \quad (10)$$

Argue that this is only satisfied when one of the two first arguments of G_+ is integer, implying

$$\pm\omega = \Delta + \ell + 2n \quad (n = 0, 1, 2, \dots) \quad (11)$$

- (e) Since ω is interpreted as the energy, it must be real. Argue that this gives a lower bound on m^2 , and give its value.

2. Lowest scalar energy state in AdS

Now consider AdS_{d+1} in global coordinates

$$ds^2 = -\cosh(\rho)^2 dt^2 + d\rho^2 + \sinh(\rho)^2 d\Omega_{d-1}^2 \quad (12)$$

- (a) (*Optional*) Show that, in global coordinates of AdS_{d+1} , the conformal generators take the form

$$\begin{aligned} D &= i \frac{\partial}{\partial t}, & M_{\mu\nu} &= -i \left(\Omega_\mu \frac{\partial}{\partial \Omega^\nu} - \Omega_\nu \frac{\partial}{\partial \Omega^\mu} \right), \\ P_\mu &= -ie^{-it} \left[\Omega_\mu (\partial_\rho - i \tanh \rho \partial_t) + \frac{1}{\tanh \rho} \nabla_\mu \right], \\ K_\mu &= ie^{it} \left[\Omega_\mu (-\partial_\rho - i \tanh \rho \partial_t) - \frac{1}{\tanh \rho} \nabla_\mu \right], \end{aligned}$$

where

$$\nabla_\mu = \frac{\partial}{\partial \Omega^\mu} - \Omega_\mu \Omega^\nu \frac{\partial}{\partial \Omega^\nu}$$

is the covariant derivative on the unit sphere S^{d-1} .

- (b) Impose that $K_\mu \phi = 0$, $D\phi = \Delta\phi$. Show that this restricts

$$\phi = \left(\frac{e^{-it}}{\cosh \rho} \right)^\Delta \quad (13)$$

- (c) By changing from global to conformal coordinates, make connection with the previous exercise. Show in particular that this implies that Δ 's are the same.

3. Vector fields in AdS

Consider the *Proca action* for a massive vector field in Lorentzian AdS in $D = d + 1$ dimensions:

$$I = - \int d^D x \sqrt{-g} \left(\frac{1}{4} g^{ac} g^{bd} F_{ab} F_{cd} + \frac{1}{2} m^2 g^{ab} A_a A_b \right), \quad (14)$$

where $F_{ab} \equiv \nabla_a A_b - \nabla_b A_a = \partial_a A_b - \partial_b A_a$.

In the Poincaré patch, we can write the metric as $g_{ab} = \Omega^2 \eta_{ab}$ with $\Omega = 1/z$, to get

$$I = - \int d^D x \left(\frac{1}{4} \Omega^{D-4} F_{ab} F^{ab} + \frac{1}{2} \Omega^{D-2} m^2 A_a A^a \right), \quad (15)$$

where indices are now raised with η_{ab} .

(a) How many physical degrees of freedom does a vector field contain when:

i) $m^2 > 0$, ii) $m^2 = 0$?

(b) Show that the equations of motion can be written as

$$\partial_a (z^{4-D} F^{ab}) = z^{2-D} m^2 A^b. \quad (16)$$

(c) Using the following ansatz for the asymptotic solution for the d components with $i \neq z$

$$A_i(x, z) = z^\nu J_i(x) + \mathcal{O}(z^{\nu+2}), \quad (17)$$

show that the equations of motion yield

$$\nu(\nu - 1) + (3 - d)\nu = m^2. \quad (18)$$

Consider the AdS dilatation isometry in Poincaré coordinates,

$$(x^i, z) \mapsto (x'^i, z') = (\lambda x^i, \lambda z). \quad (19)$$

Since $A = A_a dx^a$ is a bulk one-form, its boundary components transform under this diffeomorphism as

$$A'_i(x', z') = \frac{\partial x^j}{\partial x'^i} A_j(x, z) = \lambda^{-1} A_i(x, z). \quad (20)$$

We *define* Δ_J by requiring that J_i has boundary weight Δ_J under $x \mapsto \lambda x$,

$$J_i(\lambda x) = \lambda^{-\Delta_J} J_i(x), \quad (21)$$

(with the index transformation already accounted for by the one-form rule above). By applying the dilatation to $A_i(x, z) = z^\nu J_i(x) + \dots$, show that consistency implies

$$\Delta_J = \nu + 1 \quad (22)$$

(d) Now consider $m^2 = 0$. Show that we need

$$\Delta_J = d - 1 \quad \text{or} \quad \Delta_J = 1. \quad (23)$$

What are the physical interpretations of these two modes?

Hint: It might be helpful to think of which of the two modes decays the fastest, as well as the possible dual CFT operator.

(e) Since $m^2 = 0$, we have a gauge symmetry $\delta A_a = \nabla_a \alpha$, which can be fixed by imposing, for example, $A_z = 0$. Show that the z -equation of motion can be written as

$$\partial_i F^{iz} = 0 \quad (24)$$

Conclude from this that J_i can be interpreted as a conserved boundary current. In addition, check that the number of degrees of freedom of gauge vectors in D dimensions and conserved currents in d dimensions match.